# Krichever-Novikov continuous basis for plane algebraic curves 

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Krichever-Novikov continuous basis for plane algebraic curves.

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A continuous KN basis is a family of functions $\Phi(P, u)$
on an algebraic curve $V$, that is $P \in V$, numbered
by a continuous parameter $u$. It is assumed that $\Phi(P, u)$ is smooth in $u$. KN basis is characterized
by the property

$$
\Phi(P, u) \Phi(P, v)=L \Phi(P, u+v)
$$

where $L$ is a linear differential operator in $u$, not depending on the point $P$.
KN basis is the basis of Fourier-Laurent transform
on the curve $V$.
We shall start with basic definitions. Then we focus
on the construction and the properties of differential
operator $L$. We demonstrate a connection of the the
multiplicative property of the KN basis of $V$ with the addition law on the Jacobian of $V$.

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P.G.Grinevich and S.P.Novikov (Topological Charge of the real
periodic finite-gap Sine-Gordon solutions: Dedicated to the

Memory of J.K.Moser, Commun. Pure Appl. Math., $56(7), 2003,956-978)$ proposed an analog of the FourierLaurent integral transform on the Riemann surfaces. They use a continuous analog of the discrete KricheverNovikov bases, which were introduced and studied for the needs of the quantum string theory in the late 80-s. " Let us consider the following set of data: a nonsingular Riemann surface $\Gamma$ of the genus $g$ with marked
point $\infty \in \Gamma$ and selected local parameter near this point $z=k^{-1}, z(\infty)=0$. We construct a function $\psi^{0}=\psi(P, x)$ holomorphic on $\Gamma \backslash \infty$ and exponential near the infinite point:

$$
\psi(z, x)=k^{g} \exp \{k x\}\left(1+\sum_{i>0} \eta_{i}(x) k^{-i}\right)
$$

"Problem. Which multiplicative properties have the basic functions $\psi(P, x)=\psi_{x}(P)$ depending on $x$ as parameter?"
"Theorem. Let $x, y \neq 0$. There exists a differential operator $L$ in the variable $x$ of the order $g$ with
coefficients dependent on the both variables $x, y$
such that the following Almost Graded Commutative Associative Ring Structure is defined by the formula

$$
\begin{gathered}
\psi(P, x) \psi(P, y)=L \psi(P, x+y) \\
L=\partial_{x}^{g}+[\eta(x)+\eta(y)-\eta(x+y)] \partial_{x}^{g-1}+\ldots{ }^{\prime \prime}
\end{gathered}
$$

When $g=1$ one has: $L=\partial_{x}-(\zeta(x)+\zeta(y)-$ $\zeta(x+y))$.
In [4] we construct the operator $L$ for curves of higher genera. The algorithm is based on reduction of this problem to an effective description of the addition law on Jacobi variety.

## Group of covariant shifts

Def. Covariant shift

$$
\begin{aligned}
& W_{\alpha, \beta, c}(f(u)):= \\
& \quad \exp \{\pi \imath(\langle 2 u+\alpha, \beta\rangle+c)\} f(u+\alpha),
\end{aligned}
$$

where $u, \alpha, \beta \in \mathbb{C}^{g} ; c \in \mathbb{C}, \imath^{2}=-1 \quad$ and $\langle\cdot, \cdot\rangle$ is Euclidean scalar product.

Group of covariant shifts $:=S$

$$
\begin{aligned}
& W_{\alpha_{2}, \beta_{2}, c_{2}} W_{\alpha_{1}, \beta_{1}, c_{1}}= \\
& \quad W_{\alpha_{1}+\alpha_{2}, \beta_{1}+\beta_{2}, c_{1}+c_{2}+\left\langle\beta_{1}, \alpha_{2}\right\rangle-\left\langle\alpha_{1}, \beta_{2}\right\rangle}
\end{aligned}
$$

Representations of lattices
We use representations
$\mathbb{Z}^{g} \times \mathbb{Z}^{g} \rightarrow S$
defined by the formula

$$
\left(n, n^{\prime}\right) \mapsto(\alpha, \beta, c)=\left(\left(n, n^{\prime}\right) \Omega, \phi\left(n, n^{\prime}\right)\right)
$$

Where:
(1) $\quad \Omega \in \operatorname{Sp}(2 g, \mathbb{C})$ :

$$
\Omega^{t} J \Omega=J, \quad J=\left(\begin{array}{cc}
0 & 1_{g} \\
-1_{g} & 0
\end{array}\right) .
$$

(2) $\phi: \mathbb{Z}^{2 g} \rightarrow \mathbb{Z}_{2} \quad$ is an Arf function:
for all $Q_{1}$ and $Q_{2}$ in $\mathbb{Z}^{2 g} \quad$ Arf identity holds

$$
\phi\left(Q_{1}+Q_{2}\right)=\phi\left(Q_{1}\right)+\phi\left(Q_{2}\right)+Q_{1} J Q_{2}^{t} \quad \bmod 2 ;
$$

To define an Arf function $\phi$ fix $\left(\ell, \ell^{\prime}\right) \in \mathbb{Z}^{2 g}$ then
$\phi\left(n, n^{\prime}\right)=\left\langle n+\ell, n^{\prime}+\ell^{\prime}\right\rangle-\left\langle\ell, \ell^{\prime}\right\rangle \bmod 2$.

## Explicit formula of a representation

Write $\Omega$ in block form

$$
\Omega=\left(\begin{array}{ll}
\Omega_{1,1} & \Omega_{1,2} \\
\Omega_{2,1} & \Omega_{2,2}
\end{array}\right) ;
$$

then we have a representation:

$$
W_{\Omega}^{\ell, \ell^{\prime}}\left(n, n^{\prime}\right):=W_{\alpha, \beta, c},
$$

where

$$
\begin{aligned}
\alpha & =\left(n \Omega_{1,1}+n^{\prime} \Omega_{2,1}\right) \\
\beta & =\left(n \Omega_{1,2}+n^{\prime} \Omega_{2,2}\right) \\
c & =\left\langle n+\ell, n^{\prime}+\ell^{\prime}\right\rangle-\left\langle\ell, \ell^{\prime}\right\rangle
\end{aligned}
$$

Construction of Sigma-function
Let $\Omega \in \operatorname{Sp}(2 g, \mathbb{C}) \quad$ and $\quad\left|\Omega_{1,1}\right| \neq 0$.
Set

$$
G_{\Omega}(u)=\exp \left\{-\frac{\pi \imath}{2} u \varkappa u^{t}\right\}, \quad \varkappa=\Omega_{1,1}^{-1} \Omega_{1,2} .
$$

Def. $\sigma\left(u, \Omega ; \ell, \ell^{\prime}\right):=\sum_{\left(n, n^{\prime}\right) \in \mathbb{Z}^{2 g}} W_{\Omega}^{\ell, \ell^{\prime}}\left(n, n^{\prime}\right) G_{\Omega}(u)$
Theorem. $\sigma\left(u, \Omega ; \ell, \ell^{\prime}\right)$ is entire function of $u \in \mathbb{C}^{g}$ iff $\operatorname{Im} \tau$ is positive definite.
where
$\tau=\Omega_{2,1} \Omega_{1,1}^{-1}, \quad \Rightarrow \quad \Omega_{2,2}=\tau \Omega_{1,1} \varkappa+\left(\Omega_{1,1}^{t}\right)^{-1}$

## Families of Sigma-functions

Theorem. Fix $\Omega$ and ( $\ell, \ell^{\prime}$ ).

$$
\begin{aligned}
& \text { If } \quad W_{\Omega}^{\ell, \ell^{\prime}}\left(k, k^{\prime}\right) F(u)=F(u), \quad \forall\left(k, k^{\prime}\right) \in \mathbb{Z}^{2 g} \\
& \text { then } \quad F(u)=\text { cont } \cdot \sigma\left(u, \Omega ; \ell, \ell^{\prime}\right) .
\end{aligned}
$$

Fix a map

$$
\Omega: \mathbb{C}^{q} \rightarrow \operatorname{Sp}(2 g, \mathbb{C}), \quad\left|\Omega_{1,1}(\lambda)\right| \neq 0
$$

then we have
Family of Sigma-Functions:

$$
\sigma\left(u, \lambda ; \ell, \ell^{\prime}\right):=\frac{\sigma\left(u, \Omega(\lambda) ; \ell, \ell^{\prime}\right)}{\sqrt{\left|\Omega_{1,1}(\lambda)\right|}}, \quad \lambda \in \mathbb{C}^{q}
$$

## The Heat operators

Fix an arbitrary smooth vector field

$$
L=\sum_{j=1}^{q} v_{j}(\lambda) \frac{\partial}{\partial \lambda_{j}} .
$$

Introduce the second order operator

$$
H_{\Omega}=\sum_{r, s=1}^{g}\left(\alpha_{r, s} \partial_{r, s}+2 \alpha_{g+r, s} u_{r} \partial_{s}+\alpha_{g+r, g+s} u_{r} u_{s}\right)
$$

where $\quad\left(\alpha_{r, s}\right)=L\left(\Omega^{t}\right) J \Omega \quad$ and $\quad \partial_{r, s}=\frac{\partial^{2}}{\partial u_{r} \partial u_{s}}$

$$
\begin{aligned}
\delta_{\Omega}(\lambda) & =\frac{1}{2} \operatorname{sk}-\operatorname{tr}\left(L\left(\Omega^{t}\right) J \Omega\right) \\
\text { sk-tr } M & :=\sum_{i=1}^{g} m_{i, g+1-i} \text { for } M=\left(m_{i, j}\right)
\end{aligned}
$$

Lemma. For any constant $K \in \operatorname{Sp}(2 g, \mathbb{C})$

$$
H_{\Omega}=H_{K \Omega}, \quad \delta_{\Omega}(\lambda)=\delta_{K \Omega}(\lambda) .
$$

Lemma.

$$
\left(2 L+H_{\Omega}+\delta_{\Omega}(\lambda)\right) \frac{G_{\Omega}(u)}{\sqrt{\left|\Omega_{1,1}(\lambda)\right|}}=0 .
$$

Lemma. For all $\quad\left(k, k^{\prime}\right) \in \mathbb{Z}^{2 g}$ the covariant shift $\quad W_{\Omega(\lambda)}^{\ell, \ell^{\prime}}\left(k, k^{\prime}\right)$
and the heat operator $2 L+H_{\Omega}+\delta_{\Omega}(\lambda)$ commute as the operators on the space of smooth
functions of $u$ and $\lambda$.

Theorem. The family $\sigma\left(u, \lambda ; \ell, \ell^{\prime}\right)$ solves the equation

$$
\left(2 L+H_{\Omega}+\delta_{\Omega}(\lambda)\right) \sigma\left(u, \lambda ; \ell, \ell^{\prime}\right)=0
$$

Example. $q=1, g=1, \quad L=\partial_{\lambda}$

$$
\Omega(\lambda)=\left(\begin{array}{cc}
\omega & \omega \varkappa \\
\tau \omega & \tau \omega \varkappa+1 / \omega
\end{array}\right)
$$

$\omega \neq 0, \operatorname{Im} \tau>0, \varkappa$ are smooth functions in $\lambda$.

$$
\begin{aligned}
& H_{\Omega}=a_{1,1} \partial_{u}^{2}+2 a_{1,2} u \partial_{u}+a_{2,2} u^{2} \\
& \delta_{\Omega}(\lambda)=a_{1,2}
\end{aligned}
$$

where

$$
\begin{aligned}
& a_{1,1}=-\omega^{2} \partial_{\lambda} \tau, \\
& a_{1,2}=\frac{\partial_{\lambda} \omega}{\omega}-\omega^{2} \varkappa \partial_{\lambda} \tau, \\
& a_{2,2}=\partial_{\lambda} \varkappa-\omega^{2} \varkappa^{2} \partial_{\lambda} \tau+2 \varkappa \frac{\partial_{\lambda} \omega}{\omega}
\end{aligned}
$$

Example. $q=1, g$ is arbitrary

$$
\begin{aligned}
& L=\partial_{\lambda} \\
& \Omega(\lambda)=\left(\begin{array}{cc}
\omega & \omega \varkappa \\
\tau \omega & \tau \omega \varkappa+\left(\omega^{t}\right)^{-1}
\end{array}\right)
\end{aligned}
$$

where $|\omega| \neq 0, \tau^{t}=\tau, \operatorname{Im} \tau$ is positive definite, $\varkappa^{t}=\varkappa$ are smooth $(g \times g)$-matrix functions in $\lambda$.

$$
\begin{aligned}
& H_{\Omega}=\left(\partial_{u}\right)^{t} A_{1,1} \partial_{u}+2 u^{t} A_{2,1} \partial_{u}+u^{t} A_{2,2} u \\
& \delta_{\Omega}(\lambda)=\operatorname{tr} A_{2,1}
\end{aligned}
$$

where

$$
\begin{aligned}
& A_{1,1}=-\omega^{t}\left(\partial_{\lambda} \tau\right) \omega \\
& A_{2,1}=\omega^{-1} \partial_{\lambda} \omega+\varkappa A_{1,1}, \\
& A_{2,2}=\partial_{\lambda} \varkappa+A_{2,1} \varkappa+\varkappa A_{2,1}^{t}-\varkappa A_{1,1} \varkappa .
\end{aligned}
$$

## Abelian Sigma-functions and Heat Equations

in Non-holonomic frame.
Let $s>n>1, \quad \operatorname{gcd}(n, s)=1$.

$$
f(x, y, \lambda)=y^{n}-x^{s}-\sum_{j=0}^{n-2} \sum_{i=0}^{s-2} \lambda_{n s-i n-j s} x^{i} y^{j}
$$

the number $m:=\#\left\{\lambda_{k} \mid k<0\right\}$ is called modality.
Set: $\quad \lambda_{k}=0, \quad k<0$.
Def. The family of $(n, s)$-curves:
$V=V_{\lambda}=\left\{(x, y) \in \mathbb{C}^{2}, \lambda \in \mathbb{C}^{2 g-m} \mid f(x, y, \lambda)=0\right\}$
genus of a generic curve $V_{\lambda}$ is $g=\frac{(n-1)(s-1)}{2}$.
Example. Hyperelliptic curves are ( $2,2 g+1$ )curves, $m=0$ :

$$
f(x, y, \lambda)=y^{2}-x^{2 g+1}-\sum_{i=0}^{2 g-1} \lambda_{2(2 g-i+1)} x^{i}
$$

The principal Arf function of $(n, s)$-curve
Def. Weierstrass sequence $\left(w_{1}, \ldots, w_{g}\right)$ is the ordered set $\mathbb{N} \backslash \mathcal{M}$, where $\mathcal{M}=\{a n+b s\}, a, b \in$ $\mathbb{N} \cup 0$.
Assign $w(\xi)=\sum_{i} \xi^{w_{i}} \quad$ then

$$
w(\xi)=\frac{1}{1-\xi}-\frac{1-\xi^{n s}}{\left(1-\xi^{n}\right)\left(1-\xi^{s}\right)} .
$$

Now $g=w(1)=\frac{(n-1)(s-1)}{2}$.
Let us define ( $\pi_{1}, \ldots, \pi_{g}$ ) by the formula

$$
\begin{aligned}
& \quad \pi_{k}=w_{g-k+1}-(g-k), \quad k=1, \ldots, g . \\
& \left(w_{g}=2 g-1 \text { and } \pi_{1}=g .\right)
\end{aligned}
$$

Def. The principal Arf function is defined by

$$
\ell=(1, \ldots, 1) \quad \text { and } \quad \ell^{\prime}=\left(\pi_{1}, \ldots, \pi_{g}\right) .
$$

Let us fix this value of $\left(\ell, \ell^{\prime}\right)$ for the rest of the talk.
Example. For hyperelliptic family $(n, s)=$ $(2,2 g+1)$

$$
\ell^{\prime}=(g, g-1, \ldots, 1) .
$$

Meromorphic Abelian integrals
We do the hyperelliptic case for simplicity. Set:

$$
\begin{aligned}
& \mathrm{d} u_{j}(x, y)=\frac{x^{j-1} \mathrm{~d} x}{2 y}, \quad j=1, \ldots, g \\
& \mathrm{~d} r_{j}(x, y)=-\left(x \partial_{x} \pi_{+}\left(\frac{f(x, y, \lambda)}{x^{2 j}}\right)\right) \mathrm{d} u_{j}(x, y),
\end{aligned}
$$

where $\pi_{+}(\cdot)$ truncates the negative powers of $x$.

Def. Universal Abelian cover of $V$ is the space W
of pairs $((x, y) ;[\gamma])$, where $(x, y) \in V ;[\gamma]$ is an equivalence class of paths from $\infty \in V$ to $(x, y)$.
We say that $\gamma_{1}$ and $\gamma_{2}$ are in $[\gamma]$ if the contour $\gamma_{1} \circ \gamma_{2}^{-1}$ is homologous to zero.
Def. Abelian maps: $A: W \rightarrow \mathbb{C}^{g}, A^{*}: W \rightarrow$ $\mathbb{C}^{g}$,

$$
\begin{aligned}
A_{j}(x, y ;[\gamma]) & =\int_{\gamma} \mathrm{d} u_{j}(x, y), \\
A_{j}^{*}(x, y ;[\gamma]) & =\int_{\gamma} \mathrm{d} r_{j}(x, y) .
\end{aligned}
$$

The period map $\quad \Omega: \mathbb{C}^{2 g-m} \rightarrow \operatorname{Sp}(2 g, \mathbb{C})$
Let the contours $\gamma_{1}, \ldots, \gamma_{2 g}$ give a basis in $H_{1}(V, \mathbb{Z})$, such that the intersection matrix is $J$ :
$J_{a, b}=\gamma_{a} \circ \gamma_{b}=\operatorname{sign}(b-a) \delta_{g,|b-a|}, \quad a, b=1, \ldots, 2 g$.
Set for $i, j=1, \ldots, g$,
$\omega_{i, j}=\frac{1}{2} \oint_{\gamma_{j}} \mathrm{~d} u_{i}(x, y), \quad \omega_{i, j}^{\prime}=\frac{1}{2} \oint_{\gamma_{g+j}} \mathrm{~d} u_{i}(x, y)$,
$\eta_{i, j}=-\frac{1}{2} \oint_{\gamma_{j}} \mathrm{~d} r_{i}(x, y), \quad \eta_{i, j}^{\prime}=-\frac{1}{2} \oint_{\gamma_{g+j}} \mathrm{~d} r_{i}(x, y)$,
then for $\quad \Omega(\lambda)=\left(\begin{array}{cc}\omega & \eta \\ \omega^{\prime} & \eta^{\prime}\end{array}\right)$, where $\omega=\left(\omega_{i, j}\right)$, etc.,
we have Legendre relation

$$
\Omega(\lambda) J \Omega(\lambda)^{t}=\frac{\pi \imath}{2} J
$$

Note: For nonsingular $V_{\lambda}$ the choice of basis contours provides $|\omega| \neq 0$ and positive definiteness of $\operatorname{Im} \omega^{\prime} \omega^{-1}$. Note: $K=\ell \omega+\ell^{\prime} \omega^{\prime}$ is the vector of Riemann constants.

Frame tangent to Discriminant and Heat operators Denote by $\Delta(\lambda)$ the discriminant of $f(x, y, \lambda)$ :

$$
\Delta(\lambda)=0 \Leftrightarrow \exists(x, y) f=f_{x}=f_{y}=0
$$

Consider the space $\mathcal{T}$ of polynomial vector fields tangent to $\left\{\lambda \in \mathbb{C}^{2 g-m} \mid \Delta(\lambda)=0\right\}$. $L \in \mathcal{T}$ implies

$$
L \Delta(\lambda)=\varphi(\lambda) \Delta(\lambda), \quad \varphi(\lambda) \in \mathbb{C}[\lambda] .
$$

$\mathcal{T}$ has the basis $\left\{L_{1}, \ldots, L_{2 g}\right\}$ over $\mathbb{C}[\lambda]$. The basis gives a non-holonomic frame which defines a nontrivial polynomial Lie algebra (the structure is described in [1]).

We use the basis vector fields $\left\{L_{1}, \ldots, L_{2 g}\right\}$ and the period map $\Omega(\lambda)$ to construct $2 g$ heat operators.

Examples of basis fields $L_{i}$

$$
\begin{gathered}
g=1 . \text { We have } \Delta(\lambda)=4 \lambda_{4}^{3}+27 \lambda_{6}^{2} \\
L_{0}=4 \lambda_{4} \partial_{4}+6 \lambda_{6} \partial_{6} \\
L_{2}=6 \lambda_{6} \partial_{4}-\frac{4}{3} \lambda_{4}^{2} \partial_{6} .
\end{gathered}
$$

Here $\partial_{k}=\frac{\partial}{\partial \lambda_{k}}, \quad \operatorname{deg} \lambda_{k}=k$. Then $\quad \operatorname{deg} L_{j}=$ $j$.
$g=2$. The symmetric matrix $T$ transforms the standard fields $\partial_{4}, \partial_{6}, \partial_{8}, \partial_{10}$ to the basis fields
$L_{0}, L_{2}, L_{4}, L_{6}$
$T=\left(\begin{array}{cccc}4 \lambda_{4} & 6 \lambda_{6} & 8 \lambda_{8} & 10 \lambda_{10} \\ * & \frac{40 \lambda_{8}-12 \lambda_{4}^{2}}{5} & \frac{50 \lambda_{10}-8 \lambda_{4} \lambda_{6}}{5} & -\frac{4 \lambda_{4} \lambda_{8}}{5} \\ * & * & \frac{20 \lambda_{4} \lambda_{8}-12 \lambda_{6}^{2}}{5} & \frac{30 \lambda_{4} \lambda_{10}-6 \lambda_{6} \lambda_{8}}{5} \\ * & * & * & \frac{4 \lambda_{6} \lambda_{10}-8 \lambda_{8}^{2}}{5}\end{array}\right)$
Note: $\quad \Delta(\lambda)=|T|$. The matrix $T$ plays an important role
in Singularity Theory as the convolution matrix.
$\frac{\text { Examples of operators }}{\partial} H_{i}$
Here $D_{i}=\frac{\partial}{\partial u_{i}} ; \quad \operatorname{deg} u_{i}=-i \quad \operatorname{deg} H_{i}=i$. $g=1$.

$$
\begin{aligned}
& H_{0}=u_{1} D_{1}-1 \\
& 6 H_{2}=3 D_{1}^{2}-\lambda_{4} u_{1}^{2} \\
& g=2 . \\
& H_{0}=u_{1} D_{1}+3 u_{3} D_{3}-3 \\
& 10 H_{2}=5 D_{1}^{2}+10 u_{1} D_{3}-8 \lambda_{4} u_{3} D_{1}- \\
&-3 \lambda_{4} u_{1}^{2}+\left(15 \lambda_{8}-4 \lambda_{4}^{2}\right) u_{3}^{2}
\end{aligned}
$$

$5 H_{4}=5 D_{1} D_{3}+5 \lambda_{4} u_{3} D_{3}-6 \lambda_{6} u_{3} D_{1}-5 \lambda_{4}-$

$$
-\lambda_{6} u_{1}^{2}+5 \lambda_{8} u_{1} u_{3}+3\left(5 \lambda_{10}-\lambda_{4} \lambda_{6}\right) u_{3}^{2}
$$

$10 H_{6}=5 D_{3}^{2}-6 \lambda_{8} u_{3} D_{1}-5 \lambda_{6}-$

$$
-\lambda_{8} u_{1}^{2}+20 \lambda_{10} u_{1} u_{3}-3 \lambda_{4} \lambda_{8} u_{3}^{2}
$$

Abelian Sigma-function.
Our construction gives the following result.
Theorem. The heat equations

$$
L_{i} \sigma(u, \lambda)=H_{i} \sigma(u, \lambda)
$$

for $i \in\{n k+j s\}, 0 \leqslant j<n-1,0 \leqslant k<s-1$, uniquely define the Abelian $\sigma$-function of $(n, s)$ curve.
(1) It has the translation property

$$
\begin{aligned}
& \sigma(u+A[\chi])=\sigma(u) \exp \left\{-\left\langle A^{*}[\chi], u+\frac{1}{2} A[\chi]\right\rangle+\right. \\
&\left.+\pi \imath\left(\left\langle k+\ell, k^{\prime}+\ell^{\prime}\right\rangle-\left\langle\ell, \ell^{\prime}\right\rangle\right)\right\},
\end{aligned}
$$

where $[\chi]=\sum_{j=1}^{g}\left(k_{j} \gamma_{j}+k_{j}^{\prime} \gamma_{g+j}\right)$.
(2) It is an entire function on $\mathbb{C}^{g} \times \mathbb{C}^{2 g-m}$.

Its power series in $u$ and $\lambda$ has rational coefficients.
(3) The grading $\operatorname{deg} x=n$, $\operatorname{deg} y=s$ and $\operatorname{deg} \lambda_{k}=k$ gives $\operatorname{deg} f(x, y, \lambda)=n s$ and the grading of $u$ s.t.

$$
\operatorname{deg} \sigma(u, \lambda)=-\sum_{j=1}^{g} \ell_{j}^{\prime}=-\frac{\left(n^{2}-1\right)\left(s^{2}-1\right)}{24}
$$

## Krichever-Novikov continuous basis

Def.

$$
\psi(x, y ;[\gamma]):=\exp \left\{-\int_{[\gamma]}\left\langle A^{*}\left(\left(x^{\prime}, y^{\prime}\right),\left[\gamma^{\prime}\right]\right), \mathrm{d} A\left(x^{\prime}, y^{\prime}\right)\right\rangle\right\}
$$

$\psi(x, y ;[\gamma])$ is the unique entire function $W \rightarrow \mathbb{C}$ with:
(1) Single essentially singular point $\infty \in V$

$$
\psi \sim \xi^{g}(1+O(\xi)) .
$$

(2) No zeros and poles in $V \backslash \infty$.
$\xi$ is local parameter at $\infty, \operatorname{deg} \xi=-1$.
Def.

$$
\Psi(u,(x, y)):=\frac{\sigma(A(x, y ;[\gamma])-u)}{\psi(x, y ;[\gamma]) \sigma(u)} \exp ^{-\left\langle A^{*}(x, y ;[\gamma]), u\right\rangle}
$$

$\Psi$ is single-valued function $\mathbb{C}^{g} \times V \rightarrow \mathbb{C}$.
If $g=1, \quad \psi\left(u,\left(\wp(\xi), \wp^{\prime}(\xi)\right)\right)=\frac{\sigma(\xi-u)}{\sigma(\xi) \sigma(u)} \exp \{u \zeta(\xi)\}$.
gives a solution of Lamè equation
$\partial_{u}^{2} \Psi(u,(x, y))-2 \wp(u) \psi(u,(x, y))=x \psi(u,(x, y))$

Fix $u \in \mathbb{C}^{g}$. Then $\Psi(u,(x, y))$ is the unique single-valued function on $V$ with:
(1) $g$ zeros on $V$ at $A^{-1}(u)$.
(2) Single essentially singular point $\infty \in V$

$$
\Psi \sim \xi^{-g} \exp \left\{p\left(\xi^{-1} ; u, \lambda\right)\right\}(1+O(\xi)),
$$

where $p(t ; u, \lambda)=p_{1}(u, \lambda) t+\cdots+p_{2 g-1}(u, \lambda) t^{2 g-1}$ is fixed by the choice of $f(x, y, \lambda)$.
$p_{k}(u, \lambda)$ is homogeneous polynomial deg $p_{k}(u, \lambda)=$ $-k$. In general case

$$
\begin{aligned}
& p_{1}\left(u_{1}, 0, \ldots, 0, \lambda\right)=u_{1}, \\
& p_{j}\left(u_{1}, 0, \ldots, 0, \lambda\right)=0, \quad j>1 .
\end{aligned}
$$

Note: $\Psi(u,(x, y))$ is the Baker-Akhiezer Function corresponding to the degenerate set of Krichever data.

Example. For hyperelliptic curves

$$
\begin{aligned}
& p(t ; u, 0)=\sum_{i=1}^{g} u_{2 i-1} t^{2 i-1} \\
& g=1, \quad p(t ; u, \lambda)=u_{1} t \\
& g=2, \quad p(t ; u, \lambda)=u_{1} t+u_{3} t^{3} \\
& g=3, \quad p(t ; u, \lambda)=\left(u_{1}+\frac{1}{2} \lambda_{4} u_{5}\right) t+u_{3} t^{3}+u_{5} t^{5}
\end{aligned}
$$

Also, the equation

$$
\partial_{u_{1}}^{2} \Phi-2 \wp_{1,1}(u) \Phi=x \Phi
$$

where

$$
\wp_{1,1}(u)=-\frac{\partial^{2}}{\partial u_{1}^{2}} \log \sigma(u)
$$

has solutions

$$
\begin{aligned}
& \Phi_{ \pm}=\Psi( \pm u,(x, y)) \\
& \qquad\left|\begin{array}{rr}
\partial_{u_{1}} \Phi_{+} & \partial_{u_{1}} \Phi_{-} \\
\Phi_{+} & \Phi_{-}
\end{array}\right|=2 y
\end{aligned}
$$

In our notation the " $\psi(P, t)$ " of Grinevich and Novikov is $\Psi\left(t \mathrm{e}_{1},(x, y)\right)$, where $\mathrm{e}_{1}$ is the 1 -st ort in $\mathbb{C}^{g}$.

The relation defining the multiplicative structure of the base $\left\{\Psi\left(t e_{1},(x, y)\right)\right\}$ is a particular case of the relation

$$
\Psi(u,(x, y)) \Psi(v,(x, y))=\left.L \Psi(w,(x, y))\right|_{w=u+v}
$$

where $u, v \in \mathbb{C}^{g}$ and

$$
L=\sum_{j=0}^{g} a_{j}(u, v, w) \frac{\partial^{g-i}}{\partial w_{1}^{g-i}}, \quad \operatorname{deg} L=g .
$$

We define the family of functions on $V$ with parameter $w \in \mathbb{C}^{g}$

$$
G_{k}^{(w)}(x, y)=\frac{\partial_{w_{1}}^{k} \Psi(w,(x, y))}{\Psi(w,(x, y))}, \quad k=0,1, \ldots
$$

Each $G_{k}^{(w)}(x, y)$ is rational function on $V$.
It has $g+k$ poles in $\left\{k \infty, A^{-1}(w)\right\}$.
Its coefficients are Abelian functions on the Jacobi variety of $V$.

$$
\begin{aligned}
& G_{0}^{(w)}(x, y)=1, \\
& G_{1}^{(w)}(x, y)=-\left(\zeta_{1}(A(x, y ;[\gamma])-w)+\zeta_{1}(w)+\right. \\
&
\end{aligned}
$$

where $\zeta_{1}(w)=\partial_{w_{1}} \log \sigma(w)$.
For $k>1$ we have the recurrence
$G_{k+1}^{(w)}(x, y)=\partial_{w_{1}} G_{k}^{(w)}(x, y)+G_{1}^{(w)}(x, y) G_{k}^{(w)}(x, y)$.
We express $G_{k+1}^{(w)}(x, y)$ as rational functions of $(x, y)$.

Example. In the hyperelliptic case

$$
\begin{aligned}
& G_{1}^{(w)}(x, y)=\frac{1}{2} \frac{2 y+\sum_{i=1}^{g} \wp_{1,1,(g-i)}(w) x^{g-i}}{x^{g}-\sum_{i=1}^{g} \wp_{1, i}(w) x^{g-i}}, \\
& G_{2}^{(w)}(x, y)=x+2 \wp_{1,1}(w)
\end{aligned}
$$

from the recurrence we have for $k>2$

$$
\begin{aligned}
& G_{k}^{(w)}(x, y)=a_{k}+b_{k} G_{1}^{(w)}(x, y) \\
& a_{k+1}=\partial_{w_{1}} a_{k}+(x+2 \wp 1,1(w)) b_{k}, \\
& b_{k+1}=\partial_{w_{1}} b_{k}+a_{k},
\end{aligned}
$$

Clearly, $a_{k}$ and $b_{k}$ are polynomials in $x$.

$$
\begin{aligned}
\wp_{i, j}(w) & =-\frac{\partial^{2} \log \sigma(w)}{\partial w_{i} \partial w_{j}}, \\
\wp_{i, j, k}(w) & =-\frac{\partial^{3} \log \sigma(w)}{\partial w_{i} \partial w_{j} \partial w_{k}}
\end{aligned}
$$

where $i, j, k$ are any odd integers between 0 and $2 g$.

We prove that for all $u, v, w \in \mathbb{C}^{g}$

$$
\begin{gathered}
\Psi(u,(x, y)) \Psi(v,(x, y)) \Psi(-u-v,(x, y))=R_{3 g}^{(u, v)}(x, y) \\
\Psi(w,(x, y)) \Psi(-w,(x, y))=R_{2 g}^{(w)}(x, y)
\end{gathered}
$$

As function on $V$
$R_{3 g}^{(u, v)}(x, y)$ has
$3 g$-tuple pole at $\infty$ and $3 g$ zeros at
$\left\{A^{-1}(u), A^{-1}(v), A^{-1}(-u-v)\right\}$.
$R_{2 g}^{(w)}(x, y)$ has
$2 g$-tuple pole at $\infty$ and $2 g$ zeros at $\left\{A^{-1}(w), A^{-1}(-w)\right\}$.

The functions $R_{3 g}^{(u, v)}(x, y)$ and $R_{2 g}^{(w)}(x, y)$ define addition and inverse operations on Sym $^{g}(V)$.

Theorem. The operator $L$ is defined by the equality

$$
\frac{R_{3 g}^{(u, v)}(x, y)}{R_{2 g}^{(u+v)}(x, y)}=\sum_{i=0}^{g} a_{i}(u, v, u+v) G_{i}^{(u+v)}(x, y)
$$

This reduces the problem to comparing the coefficient at monomials, after cancelation of the common
denominator on both sides.

Example. For hyperelliptic curves we have

$$
\begin{aligned}
\alpha_{0}(u, v, w) & =1, \\
\alpha_{1}(u, v, w) & =-\zeta_{1}(u)-\zeta_{1}(v)+\zeta_{1}(w), \\
2 \alpha_{2}(u, v, w) & =-\wp_{1,1}(u)-\wp_{1,1}(v)-3 \wp_{1,1}(w)+ \\
& \left.+\alpha_{1}(u, v, w)^{2}\right), \quad \text { etc. }
\end{aligned}
$$

