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# Krichever-Novikov continuous basis for plane algebraic curves.

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A continuous KN basis is a family of functions  $\Phi(P, u)$ 

on an algebraic curve V, that is  $P \in V,$  numbered

by a continuous parameter u. It is assumed that  $\Phi(P, u)$  is smooth in u. KN basis is characterized

by the property

 $\Phi(P, u)\Phi(P, v) = L\Phi(P, u + v)$ 

where L is a linear differential operator in u, not depending on the point P.

KN basis is the basis of Fourier-Laurent transform

on the curve V.

We shall start with basic definitions. Then we focus

on the construction and the properties of differential

operator L. We demonstrate a connection of the the

multiplicative property of the KN basis of V with the addition law on the Jacobian of V.

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P.G.Grinevich and S.P.Novikov (Topological Charge of the real

periodic finite-gap Sine-Gordon solutions: Dedicated to the

Memory of J.K.Moser, Commun. Pure Appl. Math., 56(7), 2003, 956-978) proposed an analog of the Fourier-Laurent integral transform on the Riemann surfaces. They use a continuous analog of the discrete Krichever-Novikov bases, which were introduced and studied for the needs of the quantum string theory in the late 80-s. "Let us consider the following set of data: a nonsingular Riemann surface  $\Gamma$  of the genus g with marked

point  $\infty \in \Gamma$  and selected local parameter near this point  $z = k^{-1}$ ,  $z(\infty) = 0$ . We construct a function  $\psi^0 = \psi(P, x)$  holomorphic on  $\Gamma \setminus \infty$ and exponential near the infinite point:

$$\psi(z,x) = k^g \exp\{kx\}(1 + \sum_{i>0} \eta_i(x)k^{-i}).$$
 "

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"**Problem.** Which multiplicative properties have the basic functions  $\psi(P, x) = \psi_x(P)$  depending on x as parameter?"

"Theorem. Let  $x, y \neq 0$ . There exists a differential operator L in the variable x of the order g with

coefficients dependent on the both variables x, y

such that the following **Almost Graded Commutative Associative Ring Structure** is defined by the formula

$$\psi(P, x)\psi(P, y) = L\psi(P, x + y)$$

 $L = \partial_x^g + [\eta(x) + \eta(y) - \eta(x+y)]\partial_x^{g-1} + \dots$ 

When g = 1 one has:  $L = \partial_x - (\zeta(x) + \zeta(y) - \zeta(x+y)).$ 

In [4] we construct the operator L for curves of higher genera. The algorithm is based on reduction of this problem to an effective description of the addition law on Jacobi variety.

# Group of covariant shifts

### Def. Covariant shift

$$\begin{split} W_{\alpha,\beta,c}(f(u)) &:= \\ & \exp\{\pi i \left( \langle 2u + \alpha, \beta \rangle + c \right) \} f(u + \alpha), \\ & \text{where } u, \alpha, \beta \in \mathbb{C}^g \text{ ; } c \in \mathbb{C}, \ i^2 = -1 \quad \text{and} \\ & \langle \cdot, \cdot \rangle \text{ is Euclidean scalar product.} \end{split}$$

Group of covariant shifts := S

$$W_{\alpha_2,\beta_2,c_2}W_{\alpha_1,\beta_1,c_1} = W_{\alpha_1+\alpha_2,\ \beta_1+\beta_2,\ c_1+c_2+\langle\beta_1,\alpha_2\rangle-\langle\alpha_1,\beta_2\rangle}$$

Representations of lattices We use representations  $\mathbb{Z}^{g} \times \mathbb{Z}^{g} \to S$ defined by the formula

$$(n,n') \mapsto (\alpha,\beta,c) = ((n,n')\Omega, \phi(n,n')),$$

Where:

(1) 
$$\Omega \in \operatorname{Sp}(2g, \mathbb{C})$$
:  
 $\Omega^t J\Omega = J, \qquad J = \begin{pmatrix} 0 & 1_g \\ -1_g & 0 \end{pmatrix}.$   
(2)  $\phi : \mathbb{Z}^{2g} \to \mathbb{Z}_2$  is an Arf function:  
for all  $Q_1$  and  $Q_2$  in  $\mathbb{Z}^{2g}$  Arf identity holds  
 $\phi(Q_1 + Q_2) = \phi(Q_1) + \phi(Q_2) + Q_1 J Q_2^t \mod 2;$   
To define an Arf function  $\phi$  fix  
 $(\ell, \ell') \in \mathbb{Z}^{2g}$  then

 $\phi(n,n') = \langle n+\ell, n'+\ell' \rangle - \langle \ell, \ell' \rangle \mod 2.$ 

# Explicit formula of a representation

Write  $\boldsymbol{\Omega}$  in block form

$$\Omega = \begin{pmatrix} \Omega_{1,1} & \Omega_{1,2} \\ \Omega_{2,1} & \Omega_{2,2} \end{pmatrix};$$

then we have a representation:

$$W_{\Omega}^{\ell,\ell'}(n,n') := W_{\alpha,\beta,c},$$

where

$$\alpha = (n\Omega_{1,1} + n'\Omega_{2,1})$$
  

$$\beta = (n\Omega_{1,2} + n'\Omega_{2,2})$$
  

$$c = \langle n + \ell, n' + \ell' \rangle - \langle \ell, \ell' \rangle$$

$$G_{\Omega}(u) = \exp\left\{-\frac{\pi i}{2} \, u \varkappa u^t\right\}, \qquad \varkappa = \Omega_{1,1}^{-1} \Omega_{1,2}.$$

**Def.** 
$$\sigma(u,\Omega;\ell,\ell') := \sum_{(n,n')\in\mathbb{Z}^{2g}} W_{\Omega}^{\ell,\ell'}(n,n') G_{\Omega}(u)$$

**Theorem**.  $\sigma(u, \Omega; \ell, \ell')$  is entire function of  $u \in \mathbb{C}^g$  iff Im  $\tau$  is positive definite.

where

 $\tau = \Omega_{2,1} \Omega_{1,1}^{-1}, \quad \Rightarrow \quad \Omega_{2,2} = \tau \Omega_{1,1} \varkappa + (\Omega_{1,1}^t)^{-1}$ 

# Families of Sigma-functions

**Theorem**. Fix  $\Omega$  and  $(\ell, \ell')$ .

If 
$$W_{\Omega}^{\ell,\ell'}(k,k') F(u) = F(u), \quad \forall (k,k') \in \mathbb{Z}^{2g}$$
  
then  $F(u) = \text{const} \cdot \sigma(u,\Omega;\ell,\ell').$ 

Fix a map

$$\Omega: \mathbb{C}^q \to \mathsf{Sp}(2g, \mathbb{C}), \qquad |\Omega_{1,1}(\lambda)| \neq 0$$

then we have

### Family of Sigma-Functions:

$$\sigma(u,\lambda;\ell,\ell') := \frac{\sigma(u,\Omega(\lambda);\ell,\ell')}{\sqrt{|\Omega_{1,1}(\lambda)|}}, \quad \lambda \in \mathbb{C}^q.$$

The Heat operators

Fix an arbitrary smooth vector field

$$L = \sum_{j=1}^{q} v_j(\lambda) \frac{\partial}{\partial \lambda_j}.$$

Introduce the second order operator

$$H_{\Omega} = \sum_{r,s=1}^{g} (\alpha_{r,s}\partial_{r,s} + 2\alpha_{g+r,s}u_{r}\partial_{s} + \alpha_{g+r,g+s}u_{r}u_{s}),$$
  
where  $(\alpha_{r,s}) = L(\Omega^{t})J\Omega$  and  $\partial_{r,s} = \frac{\partial^{2}}{\partial u_{r}\partial u_{s}}$   
 $\delta_{\Omega}(\lambda) = \frac{1}{2}\operatorname{sk-tr}(L(\Omega^{t})J\Omega).$   
sk-tr $M := \sum_{i=1}^{g} m_{i,g+1-i}$  for  $M = (m_{i,j})$ 

**Lemma**. For any constant  $K \in \text{Sp}(2g, \mathbb{C})$  $H_{\Omega} = H_{K\Omega}, \qquad \delta_{\Omega}(\lambda) = \delta_{K\Omega}(\lambda).$  Lemma.

$$\left(2L + H_{\Omega} + \delta_{\Omega}(\lambda)\right) \frac{G_{\Omega}(u)}{\sqrt{|\Omega_{1,1}(\lambda)|}} = 0.$$

**Lemma**. For all  $(k, k') \in \mathbb{Z}^{2g}$ the covariant shift  $W_{\Omega(\lambda)}^{\ell,\ell'}(k, k')$ and the heat operator  $2L + H_{\Omega} + \delta_{\Omega}(\lambda)$ commute as the operators on the space of smooth functions of u and  $\lambda$ .

**Theorem**. The family  $\sigma(u, \lambda; \ell, \ell')$  solves the equation

$$(2L + H_{\Omega} + \delta_{\Omega}(\lambda))\sigma(u, \lambda; \ell, \ell') = 0$$

**Example**. q = 1, g = 1,  $L = \partial_{\lambda}$ 

$$\Omega(\lambda) = \begin{pmatrix} \omega & \omega \varkappa \\ \tau \omega & \tau \omega \varkappa + 1/\omega \end{pmatrix}$$

 $\omega \neq 0$ , Im $\tau > 0$ ,  $\varkappa$  are smooth functions in  $\lambda$ .

$$H_{\Omega} = a_{1,1}\partial_u^2 + 2a_{1,2}u\partial_u + a_{2,2}u^2$$
$$\delta_{\Omega}(\lambda) = a_{1,2}$$

where

$$a_{1,1} = -\omega^2 \partial_\lambda \tau,$$
  

$$a_{1,2} = \frac{\partial_\lambda \omega}{\omega} - \omega^2 \varkappa \partial_\lambda \tau,$$
  

$$a_{2,2} = \partial_\lambda \varkappa - \omega^2 \varkappa^2 \partial_\lambda \tau + 2\varkappa \frac{\partial_\lambda \omega}{\omega}$$

**Example**. q = 1, g is arbitrary

$$L = \partial_{\lambda}$$
$$\Omega(\lambda) = \begin{pmatrix} \omega & \omega \varkappa \\ \tau \omega & \tau \omega \varkappa + (\omega^{t})^{-1} \end{pmatrix}$$

where  $|\omega|\neq 0,\ \tau^t=\tau,\ {\rm Im}\tau$  is positive definite,  $\varkappa^t=\varkappa$  are smooth  $(g\times g)\text{-matrix}$  functions in  $\lambda$  .

$$H_{\Omega} = (\partial_u)^t A_{1,1} \partial_u + 2u^t A_{2,1} \partial_u + u^t A_{2,2} u$$
  
$$\delta_{\Omega}(\lambda) = \operatorname{tr} A_{2,1}$$

where

$$A_{1,1} = -\omega^t (\partial_\lambda \tau) \omega,$$
  

$$A_{2,1} = \omega^{-1} \partial_\lambda \omega + \varkappa A_{1,1},$$
  

$$A_{2,2} = \partial_\lambda \varkappa + A_{2,1} \varkappa + \varkappa A_{2,1}^t - \varkappa A_{1,1} \varkappa.$$

# Abelian Sigma-functions and Heat Equations in Non-holonomic frame.

Let s > n > 1, gcd(n, s) = 1.

$$f(x, y, \lambda) = y^n - x^s - \sum_{j=0}^{n-2} \sum_{i=0}^{s-2} \lambda_{ns-in-js} x^i y^j,$$

the number  $m := \#\{\lambda_k | k < 0\}$  is called *modality*.

Set:  $\lambda_k = 0, \quad k < 0.$ 

**Def**. The family of (n, s)-curves:

 $V = V_{\lambda} = \{(x, y) \in \mathbb{C}^2, \lambda \in \mathbb{C}^{2g-m} | f(x, y, \lambda) = 0\}$ genus of a generic curve  $V_{\lambda}$  is  $g = \frac{(n-1)(s-1)}{2}$ . **Example**. Hyperelliptic curves are (2, 2g + 1)curves, m = 0:

$$f(x, y, \lambda) = y^2 - x^{2g+1} - \sum_{i=0}^{2g-1} \lambda_{2(2g-i+1)} x^i$$

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The principal Arf function of (n, s)-curve **Def**. Weierstrass sequence  $(w_1, \ldots, w_g)$  is the ordered set  $\mathbb{N}\setminus\mathcal{M}$ , where  $\mathcal{M} = \{an + bs\}, a, b \in$  $\mathbb{N} \cup 0$ .

Assign  $w(\xi) = \sum_i \xi^{w_i}$  then

$$w(\xi) = \frac{1}{1-\xi} - \frac{1-\xi^{ns}}{(1-\xi^n)(1-\xi^s)}$$

Now  $g = w(1) = \frac{(n-1)(s-1)}{2}$ . Let us define  $(\pi_1, \dots, \pi_g)$  by the formula

$$\pi_k = w_{g-k+1} - (g-k), \qquad k = 1, \dots, g.$$

 $(w_g = 2g - 1 \text{ and } \pi_1 = g.)$ **Def**. *The principal Arf function* is defined by

$$\ell = (1, ..., 1)$$
 and  $\ell' = (\pi_1, ..., \pi_g).$ 

Let us fix this value of  $(\ell, \ell')$  for the rest of the talk.

**Example**. For hyperelliptic family (n,s) = (2,2g+1)

$$\ell'=(g,g-1,\ldots,1).$$

Meromorphic Abelian integrals We do the hyperelliptic case for simplicity. Set:

$$\mathsf{d}u_j(x,y) = \frac{x^{j-1}\mathsf{d}x}{2y}, \qquad j = 1, \dots, g,$$

$$dr_j(x,y) = -\left(x\partial_x \pi_+\left(\frac{f(x,y,\lambda)}{x^{2j}}\right)\right) du_j(x,y),$$

where  $\pi_+(\cdot)$  truncates the negative powers of x.

**Def.** Universal Abelian cover of V is the space W

of pairs  $((x, y); [\gamma])$ , where  $(x, y) \in V$ ;  $[\gamma]$  is an equivalence class of paths from  $\infty \in V$  to (x, y).

We say that  $\gamma_1$  and  $\gamma_2$  are in  $[\gamma]$  if the contour  $\gamma_1 \circ \gamma_2^{-1}$  is homologous to zero.

**Def.** Abelian maps:  $A: W \to \mathbb{C}^g$ ,  $A^*: W \to \mathbb{C}^g$ ,

$$A_j(x,y;[\gamma]) = \int_{\gamma} du_j(x,y),$$
  

$$A_j^*(x,y;[\gamma]) = \int_{\gamma} dr_j(x,y).$$
  

$$j = 1, \dots, g.$$

 $\begin{array}{ll} \hline \text{The period map} & \Omega: \mathbb{C}^{2g-m} \to \operatorname{Sp}(2g,\mathbb{C}) \\ \text{Let the contours } \gamma_1, \ldots, \gamma_{2g} \text{ give a basis in } H_1(V,\mathbb{Z}), \\ \text{such that the intersection matrix is } J: \\ J_{a,b} = \gamma_a \circ \gamma_b = \operatorname{sign}(b-a) \delta_{g,|b-a|}, \quad a,b = 1, \ldots, 2g. \\ \text{Set for } i, j = 1, \ldots, g, \end{array}$ 

$$\begin{split} \omega_{i,j} &= \frac{1}{2} \oint_{\gamma_j} \mathrm{d} u_i(x,y), \quad \omega'_{i,j} = \frac{1}{2} \oint_{\gamma_{g+j}} \mathrm{d} u_i(x,y), \\ \eta_{i,j} &= -\frac{1}{2} \oint_{\gamma_j} \mathrm{d} r_i(x,y), \quad \eta'_{i,j} = -\frac{1}{2} \oint_{\gamma_{g+j}} \mathrm{d} r_i(x,y), \end{split}$$

then for  $\Omega(\lambda) = \begin{pmatrix} \omega & \eta \\ \omega' & \eta' \end{pmatrix}$ , where  $\omega = (\omega_{i,j})$ , etc.,

we have Legendre relation

$$\Omega(\lambda)J\Omega(\lambda)^t = \frac{\pi i}{2}J$$

Note: For nonsingular  $V_{\lambda}$  the choice of basis contours provides  $|\omega| \neq 0$  and positive definiteness of  $\text{Im } \omega' \omega^{-1}$ . Note:  $K = \ell \omega + \ell' \omega'$  is the vector of *Riemann constants*. Frame tangent to Discriminant and Heat operators Denote by  $\Delta(\lambda)$  the discriminant of  $f(x, y, \lambda)$ :

 $\Delta(\lambda) = 0 \Leftrightarrow \exists (x, y) \ f = f_x = f_y = 0$ 

Consider the space  $\mathcal{T}$  of polynomial vector fields tangent to  $\{\lambda \in \mathbb{C}^{2g-m} | \Delta(\lambda) = 0\}.$  $L \in \mathcal{T}$  implies

$$L \Delta(\lambda) = \varphi(\lambda) \Delta(\lambda), \quad \varphi(\lambda) \in \mathbb{C}[\lambda].$$

 $\mathcal{T}$  has the basis  $\{L_1, \ldots, L_{2g}\}$  over  $\mathbb{C}[\lambda]$ . The basis gives a non-holonomic frame which defines a nontrivial polynomial Lie algebra (the structure is described in [1]).

We use the basis vector fields  $\{L_1, \ldots, L_{2g}\}$  and the period map  $\Omega(\lambda)$  to construct 2g heat operators. Examples of basis fields  $L_i$ 

$$g = 1. \text{ We have } \Delta(\lambda) = 4\lambda_4^3 + 27\lambda_6^2$$

$$L_0 = 4\lambda_4\partial_4 + 6\lambda_6\partial_6,$$

$$L_2 = 6\lambda_6\partial_4 - \frac{4}{3}\lambda_4^2\partial_6.$$
Here  $\partial_k = \frac{\partial}{\partial\lambda_k}$ ,  $\deg \lambda_k = k$ . Then  $\deg L_j = j$ .

g = 2. The symmetric matrix T transforms the standard fields  $\partial_4, \partial_6, \partial_8, \partial_{10}$  to the basis fields  $L_0, L_2, L_4, L_6$ 

$$T = \begin{pmatrix} 4\lambda_4 & 6\lambda_6 & 8\lambda_8 & 10\lambda_{10} \\ * & \frac{40\lambda_8 - 12\lambda_4^2}{5} & \frac{50\lambda_{10} - 8\lambda_4\lambda_6}{5} & -\frac{4\lambda_4\lambda_8}{5} \\ * & * & \frac{20\lambda_4\lambda_8 - 12\lambda_6^2}{5} & \frac{30\lambda_4\lambda_{10} - 6\lambda_6\lambda_8}{5} \\ * & * & * & \frac{4\lambda_6\lambda_{10} - 8\lambda_8^2}{5} \end{pmatrix}$$

Note:  $\Delta(\lambda) = |T|$ . The matrix T plays an important role

in Singularity Theory as the *convolution matrix*.

Examples of operators  $H_i$ Here  $D_i = \frac{\partial}{\partial u_i}$ ;  $\deg u_i = -i$   $\deg H_i = i$ . a = 1. $H_0 = u_1 D_1 - 1$  $6H_2 = 3D_1^2 - \lambda_4 u_1^2$ g = 2. $H_0 = u_1 D_1 + 3 u_3 D_3 - 3$  $10H_2 = 5D_1^2 + 10u_1D_3 - 8\lambda_4u_3D_1 -$  $-3\lambda_4 u_1^2 + (15\lambda_8 - 4\lambda_4^2)u_3^2$  $5H_4 = 5D_1D_3 + 5\lambda_4u_3D_3 - 6\lambda_6u_3D_1 - 5\lambda_4 - 5\lambda_4u_3D_3 - 6\lambda_6u_3D_1 - 5\lambda_4 - 5$  $-\lambda_6 u_1^2 + 5\lambda_8 u_1 u_3 + 3(5\lambda_{10} - \lambda_4 \lambda_6) u_3^2$  $10H_6 = 5D_3^2 - 6\lambda_8 u_3 D_1 - 5\lambda_6 -\lambda_8 u_1^2 + 20\lambda_{10}u_1u_3 - 3\lambda_4\lambda_8 u_3^2$ 

Abelian Sigma-function.

Our construction gives the following result. **Theorem**. The heat equations

$$L_i\sigma(u,\lambda) = H_i\sigma(u,\lambda),$$

for  $i \in \{nk + js\}$ ,  $0 \leq j < n - 1$ ,  $0 \leq k < s - 1$ , uniquely define the Abelian  $\sigma$ -function of (n, s)curve.

(1) It has the translation property

$$\sigma(u + A[\chi]) = \sigma(u) \exp\left\{-\left\langle A^*[\chi], u + \frac{1}{2}A[\chi]\right\rangle + \pi i \left(\langle k + \ell, k' + \ell' \rangle - \langle \ell, \ell' \rangle\right)\right\},\$$

where  $[\chi] = \sum_{j=1}^{g} (k_j \gamma_j + k'_j \gamma_{g+j}).$ (2) It is an entire function on  $\mathbb{C}^g \times \mathbb{C}^{2g-m}$ .

Its power series in u and  $\lambda$  has rational coefficients.

(3) The grading deg x = n, deg y = s and deg  $\lambda_k = k$  gives deg  $f(x, y, \lambda) = ns$  and the grading of u s.t.

deg 
$$\sigma(u, \lambda) = -\sum_{j=1}^{g} \ell'_j = -\frac{(n^2 - 1)(s^2 - 1)}{24}.$$

# Krichever-Novikov continuous basis **Def**.

$$\psi(x,y;[\gamma]) := \exp\left\{-\int_{[\gamma]} \left\langle A^*((x',y'),[\gamma']), \mathsf{d}A(x',y')\right\rangle\right\}$$

 $\psi(x,y;[\gamma])$  is the <u>unique</u> entire function  $W \to \mathbb{C}$  with:

(1) Single essentially singular point  $\infty \in V$ 

 $\psi \sim \xi^g (1 + O(\xi)).$ 

(2) No zeros and poles in  $V \setminus \infty$ .  $\xi$  is local parameter at  $\infty$ , deg  $\xi = -1$ . **Def**.

$$\Psi(u,(x,y)) := \frac{\sigma(A(x,y;[\gamma]) - u)}{\psi(x,y;[\gamma]) \,\sigma(u)} \exp^{-\langle A^*(x,y;[\gamma]),u \rangle}$$

$$\begin{split} \Psi \text{ is single-valued function } \mathbb{C}^g \times V \to \mathbb{C}. \\ \text{If } g &= 1, \quad \Psi \Big( u, (\wp(\xi), \wp'(\xi)) \Big) = \frac{\sigma(\xi - u)}{\sigma(\xi)\sigma(u)} \exp\{ u\zeta(\xi) \}. \\ \text{gives a solution of Lamè equation} \\ \partial_u^2 \Psi \Big( u, (x, y) \Big) - 2\wp(u) \Psi \Big( u, (x, y) \Big) = x \, \Psi \Big( u, (x, y) \Big) \end{split}$$

Fix  $u \in \mathbb{C}^{g}$ . Then  $\Psi(u, (x, y))$  is the **unique single-valued function** on V with: (1) g zeros on V at  $A^{-1}(u)$ . (2) Single essentially singular point  $\infty \in V$  $\Psi \sim \xi^{-g} \exp\{p(\xi^{-1}; u, \lambda)\}(1 + O(\xi)),$ where  $p(t; u, \lambda) = p_1(u, \lambda)t + \dots + p_{2g-1}(u, \lambda)t^{2g-1}$ is fixed by the choice of  $f(x, y, \lambda)$ .

 $p_k(u,\lambda)$  is homogeneous polynomial deg  $p_k(u,\lambda) = -k$ . In general case

$$p_1(u_1, 0, \dots, 0, \lambda) = u_1,$$
  
 $p_j(u_1, 0, \dots, 0, \lambda) = 0, \qquad j > 1.$ 

Note:  $\Psi(u, (x, y))$  is the *Baker-Akhiezer Function* corresponding to the degenerate set of Krichever data.

#### **Example**. For hyperelliptic curves

$$p(t; u, 0) = \sum_{i=1}^{g} u_{2i-1} t^{2i-1}.$$
  

$$g = 1, \quad p(t; u, \lambda) = u_1 t,$$
  

$$g = 2, \quad p(t; u, \lambda) = u_1 t + u_3 t^3,$$
  

$$g = 3, \quad p(t; u, \lambda) = \left(u_1 + \frac{1}{2}\lambda_4 u_5\right) t + u_3 t^3 + u_5 t^5.$$

Also, the equation

$$\partial_{u_1}^2 \Phi - 2\wp_{1,1}(u)\Phi = x \Phi,$$

where

$$\wp_{1,1}(u) = -\frac{\partial^2}{\partial u_1^2} \log \sigma(u)$$

has solutions  $\Phi_{\pm} = \Psi \left( \pm u, (x, y) \right),$   $\begin{vmatrix} \partial_{u_1} \Phi_+ & \partial_{u_1} \Phi_- \\ \Phi_+ & \Phi_- \end{vmatrix} = 2y.$ 

In our notation the 
$$(\psi(P,t))''$$
 of Grinevich  
and Novikov is  $\Psi(te_1, (x, y))$ , where  $e_1$  is the  
1-st ort in  $\mathbb{C}^g$ .

The relation defining the multiplicative structure

of the base  $\{\Psi(te_1, (x, y))\}$  is a particular case of the relation

$$\begin{split} \Psi\Bigl(u,(x,y)\Bigr)\Psi\Bigl(v,(x,y)\Bigr) &= L\Psi\Bigl(w,(x,y)\Bigr)\,\Big|_{w=u+v}, \end{split}$$
 where  $u,v\in\mathbb{C}^g$  and

$$L = \sum_{j=0}^{g} a_j(u, v, w) \frac{\partial^{g-i}}{\partial w_1^{g-i}}, \qquad \deg L = g.$$

We define the family of functions on V with parameter  $w \in \mathbb{C}^g$ 

$$G_k^{(w)}(x,y) = \frac{\partial_{w_1}^k \Psi(w,(x,y))}{\Psi(w,(x,y))}, \qquad k = 0, 1, \dots$$

Each  $G_k^{(w)}(x, y)$  is rational function on V. It has g + k poles in  $\{k \infty, A^{-1}(w)\}$ . Its coefficients are Abelian functions on the Jacobi variety of V.

$$G_0^{(w)}(x,y) = 1,$$
  

$$G_1^{(w)}(x,y) = -(\zeta_1(A(x,y;[\gamma]) - w) + \zeta_1(w) + \langle A^*(x,y;[\gamma]), e_1 \rangle),$$

where  $\zeta_1(w) = \partial_{w_1} \log \sigma(w)$ . For k > 1 we have the recurrence  $G_{k+1}^{(w)}(x,y) = \partial_{w_1} G_k^{(w)}(x,y) + G_1^{(w)}(x,y) G_k^{(w)}(x,y)$ . We express  $G_{k+1}^{(w)}(x,y)$  as rational functions of (x,y). **Example**. In the hyperelliptic case

$$G_1^{(w)}(x,y) = \frac{1}{2} \frac{2y + \sum_{i=1}^g \wp_{1,1,(g-i)}(w) x^{g-i}}{x^g - \sum_{i=1}^g \wp_{1,i}(w) x^{g-i}},$$

$$G_2^{(w)}(x,y) = x + 2\wp_{1,1}(w)$$

from the recurrence we have for k > 2

$$G_{k}^{(w)}(x,y) = a_{k} + b_{k}G_{1}^{(w)}(x,y)$$
  

$$a_{k+1} = \partial_{w_{1}}a_{k} + (x + 2\wp_{1,1}(w))b_{k},$$
  

$$b_{k+1} = \partial_{w_{1}}b_{k} + a_{k},$$

Clearly,  $a_k$  and  $b_k$  are polynomials in x.

$$\wp_{i,j}(w) = -\frac{\partial^2 \log \sigma(w)}{\partial w_i \partial w_j},$$
  
$$\wp_{i,j,k}(w) = -\frac{\partial^3 \log \sigma(w)}{\partial w_i \partial w_j \partial w_k}$$

where i, j, k are any odd integers between 0 and 2g.

We prove that for all  $u, v, w \in \mathbb{C}^g$   $\Psi(u, (x, y))\Psi(v, (x, y))\Psi(-u - v, (x, y)) = R_{3g}^{(u,v)}(x, y),$   $\Psi(w, (x, y))\Psi(-w, (x, y)) = R_{2g}^{(w)}(x, y).$ As function on V

$$egin{aligned} R^{(u,v)}_{3g}(x,y) & ext{has} \ 3g ext{-tuple pole at $\infty$ and $3g$ zeros at} \ \{A^{-1}(u),A^{-1}(v),A^{-1}(-u-v)\}. \end{aligned}$$

$$R_{2g}^{(w)}(x,y)$$
 has  
2g-tuple pole at  $\infty$  and 2g zeros at  
 $\{A^{-1}(w), A^{-1}(-w)\}.$ 

The functions  $R_{3g}^{(u,v)}(x,y)$  and  $R_{2g}^{(w)}(x,y)$  define **addition and inverse operations** on  $Sym^{g}(V)$ .

**Theorem**. The operator L is defined by the equality

$$\frac{R_{3g}^{(u,v)}(x,y)}{R_{2g}^{(u+v)}(x,y)} = \sum_{i=0}^{g} a_i(u,v,u+v)G_i^{(u+v)}(x,y).$$

This reduces the problem to comparing the coefficients at monomials, after cancelation of the common

denominator on both sides.

**Example**. For hyperelliptic curves we have

$$\begin{aligned} \alpha_0(u, v, w) &= 1, \\ \alpha_1(u, v, w) &= -\zeta_1(u) - \zeta_1(v) + \zeta_1(w), \\ 2\alpha_2(u, v, w) &= -\wp_{1,1}(u) - \wp_{1,1}(v) - 3\wp_{1,1}(w) + \\ &+ \alpha_1(u, v, w)^2 \Big), \quad \text{etc.} \end{aligned}$$